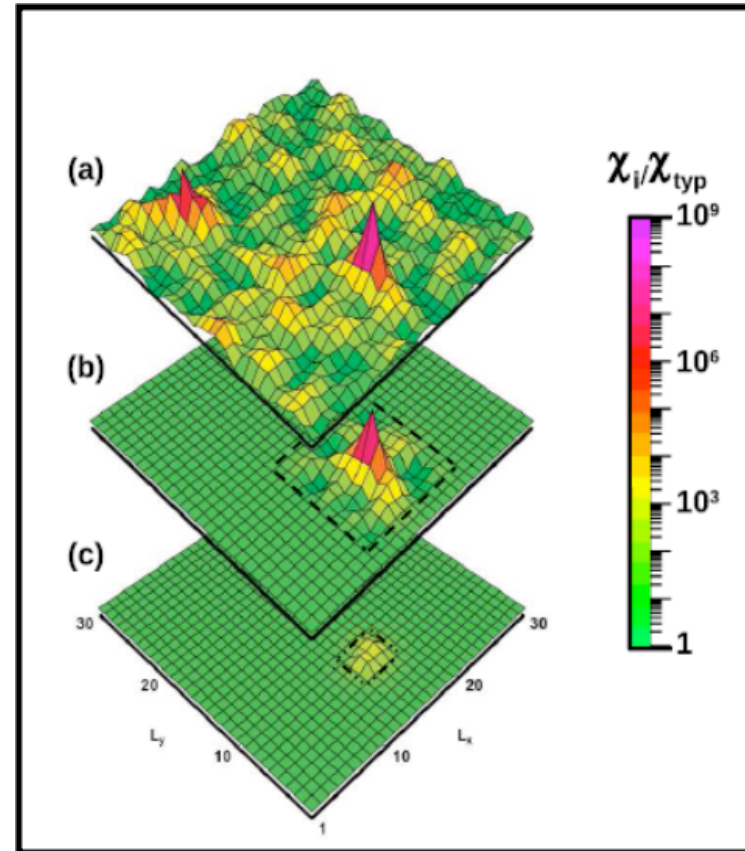
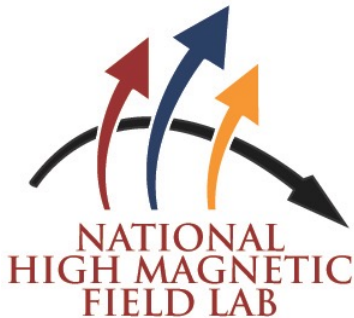


# Localization in Dynamical Mean Field Theory

## Lecture 2: Friedel Oscillations and Electronic Griffiths Phases

**Vladimir Dobrosavljevic**  
Florida State University

<http://badmetals.magnet.fsu.edu>



**Workshop “Localization in Quantum Systems”**  
Jun. 1-2, 2017, King’s College London

## Local perspective: the cavity field?

(Abou-Chacra, Thouless, Anderson (1973))

**Local effective action** (expanded to  $O(t^2)$ ): **Anderson impurity model**

$$\begin{aligned}
 S_{\text{eff}}(i) &= S_{\text{loc}}(i) - \ln \Xi(i) \\
 &= \sum_{\sigma} \int_0^{\beta} d\tau \int_0^{\beta} d\tau' c_{i,\sigma}^{\dagger}(\tau) \\
 &\quad \times [\delta(\tau - \tau') (\partial_{\tau} + \varepsilon_i - \mu) + \Delta_{i,\sigma}(\tau, \tau')] \\
 &\quad \times c_{i,\sigma}(\tau') + U \int_0^{\beta} d\tau n_{i,\uparrow}(\tau) n_{i,\downarrow}(\tau).
 \end{aligned}$$

“cavity field”?

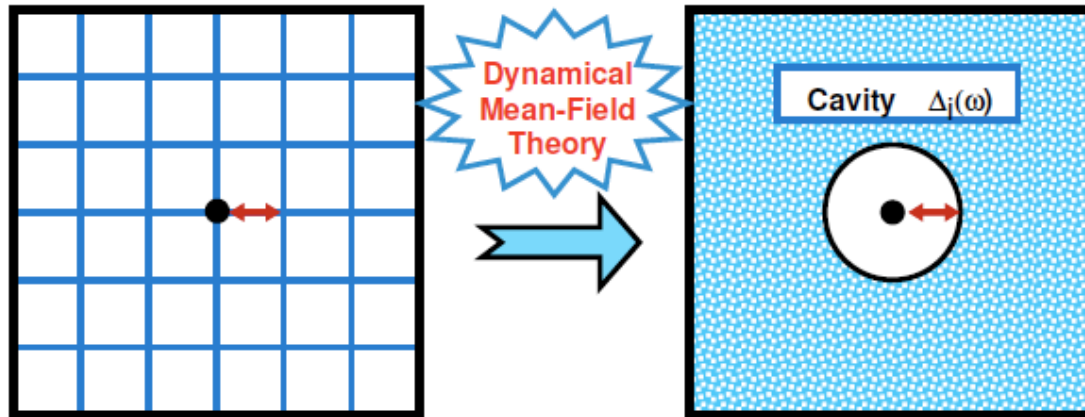
$$\Delta_i(\omega_n) = \sum_{j=1}^z t_{ij}^2 G_j^{(i)}(\omega_n)$$

fluctuates  
in energy and space

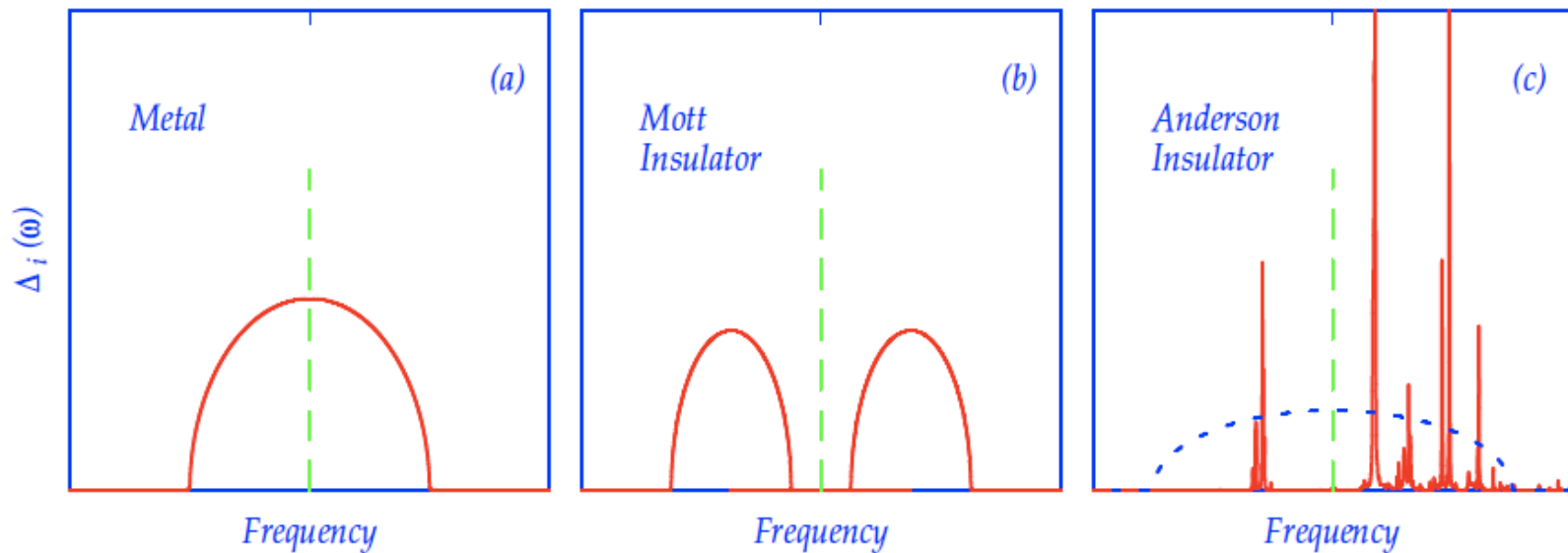
**Recursion relation:**

$$G_{cj}^{(i)(-1)}(\omega) = \omega - \varepsilon_j - \sum_{k=1}^{z-1} t_{jk}^2 G_{ck}^{(j)}(\omega)$$

# Fluctuating cavity field



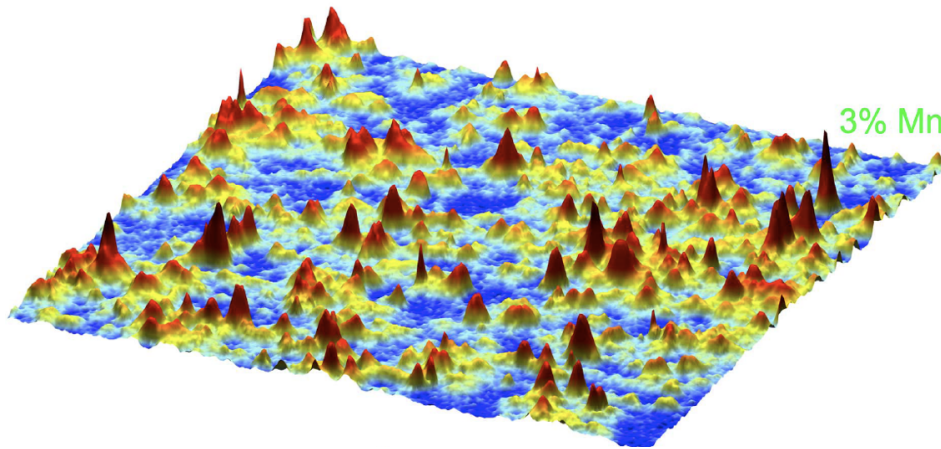
Bethe lattice simulation



# Can local spectrum recognize Anderson localization?

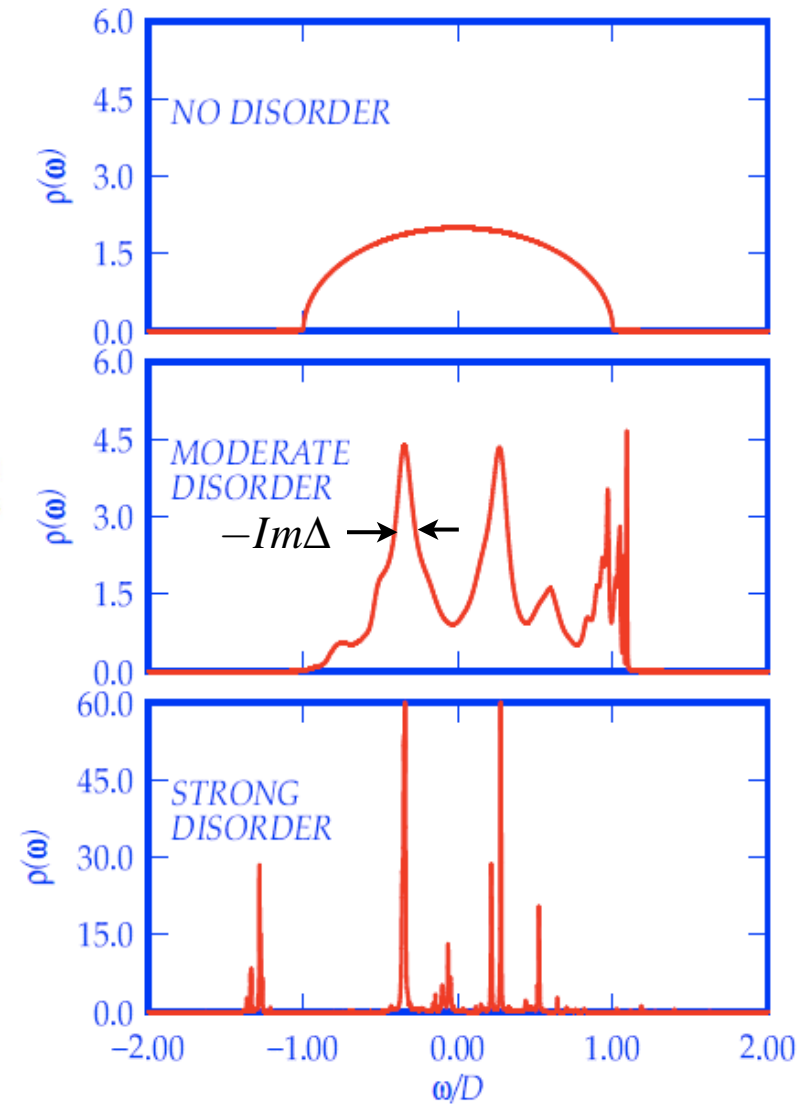
$$\rho_i(\omega) = \frac{1}{\pi} \text{Im} \frac{1}{\omega - \varepsilon_i - \Delta_i(\omega)}$$

$$= \sum_n \delta(\omega - \omega_n) |\psi_n(i)|^2$$



**Yazdani, STM experiments GaMnAs**  
(close to localization)

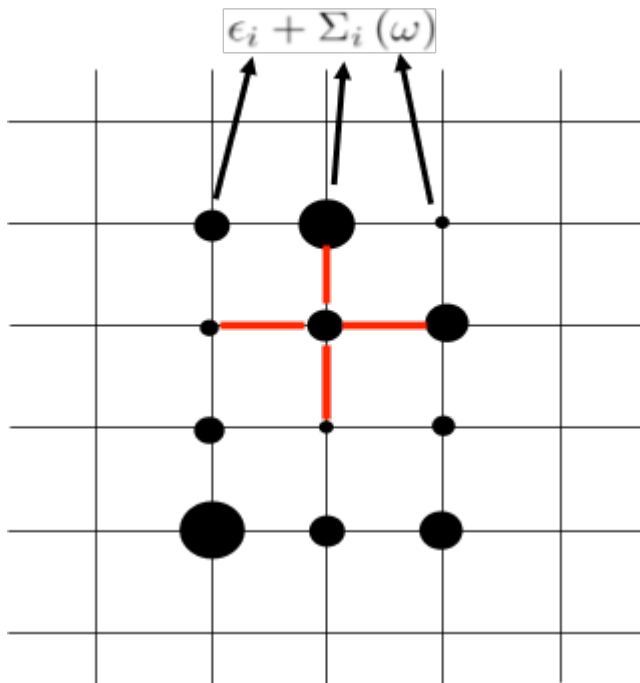
**Bethe lattice simulation**



# Disordered Mott Transitions: Quantum TAP

- Clean case ( $W=0$ ): **Mott metal-insulator transition** at  $U=U_c$ , where  $Z \rightarrow 0$  (Brinkman and Rice, 1970).
- Fermi liquid approach in which each fermion acquires a **quasi-particle renormalization** and each site-energy is **renormalized**:

Local renormalizations



$$\Sigma_i(\omega) = (1 - Z_i^{-1})\omega - \epsilon_i + \bar{\epsilon}_i/Z_i$$

**Local moment formation:**  $Z_i \rightarrow 0$

**Orbitally (site) selective Mott transition?**

“deconfinement”, “fractionalization”

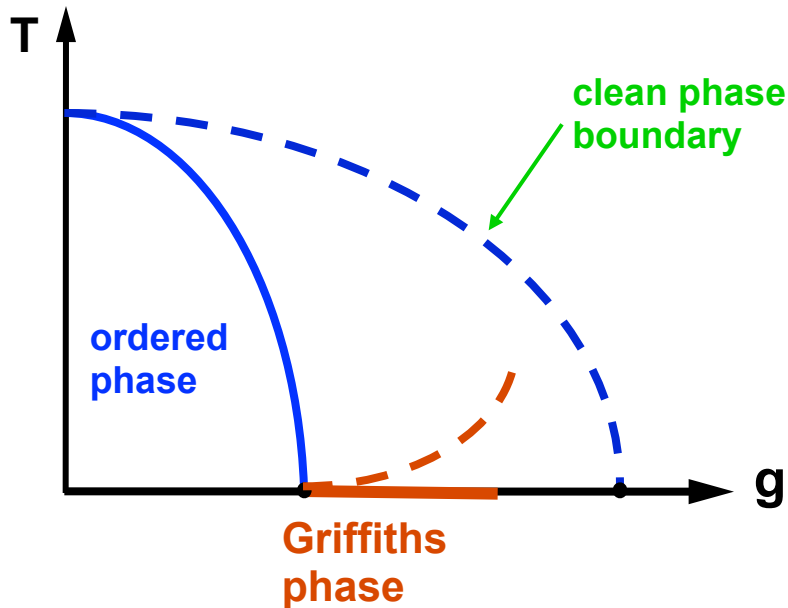
**“Kondo” THEOREM:** in any metal

$$Z_i \neq 0 \quad \rho_i \neq 0 \quad (\text{continuum spectrum})$$

(exceptions on Friday)

# Quantum Griffiths phases and IRFP (1990s)

- D. Fisher (1992): new scenario for (**insulating**) QCPs with disorder (Ising)

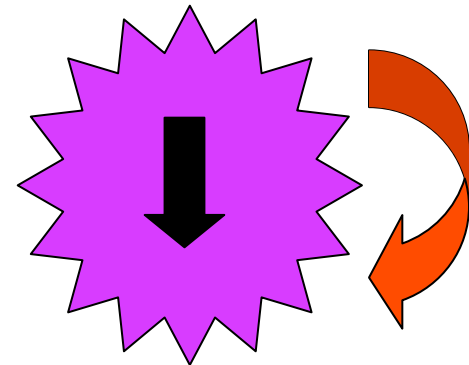


$$P(L) \sim \exp\{-\rho L^d\}$$

$$P(\Delta) \sim \Delta^{\alpha-1} ; \quad \chi \sim T^{\alpha-1}$$

$\alpha \rightarrow 0$  at QCP (**IRFP**)

Griffiths phase (**Till + Huse**):



**Rare, dilute** magnetically  
ordered cluster tunnels  
with rate  $\Delta(L) \sim \exp\{-AL^d\}$

**E.Miranda, V. Dobrosavljevic,**  
**Reports on Progress in Physics 68, 2337 (2005)**

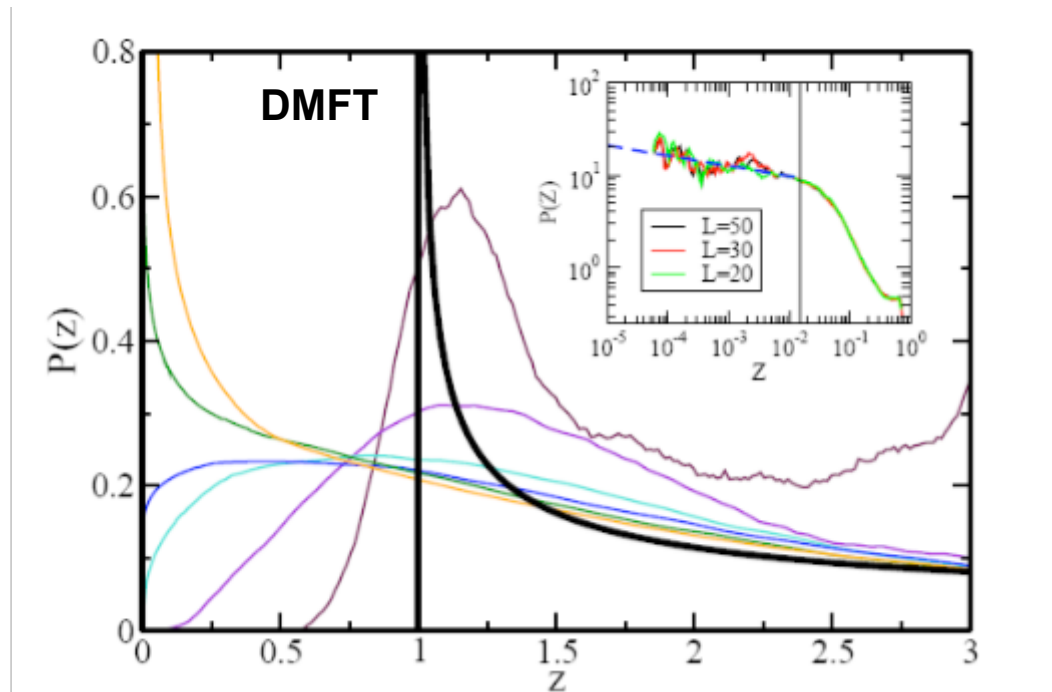
## Electronic Griffiths Phase of the $d = 2$ Mott Transition

E. C. Andrade,<sup>1,2</sup> E. Miranda,<sup>2</sup> and V. Dobrosavljević<sup>1</sup>

- In D=2, the environment of each site (“bath”) has strong **spatial fluctuations**
- **New physics: rare events** due to fluctuations and spatial correlations

Distribution  $P(Z/Z_0)$   
acquires a **low- $Z$  tail**:

$$P(Z) \propto Z^{\alpha-1}$$



# Results: Thermodynamics

- Remembering that the local Kondo temperature and  $T_{Ki} \propto Z_i$

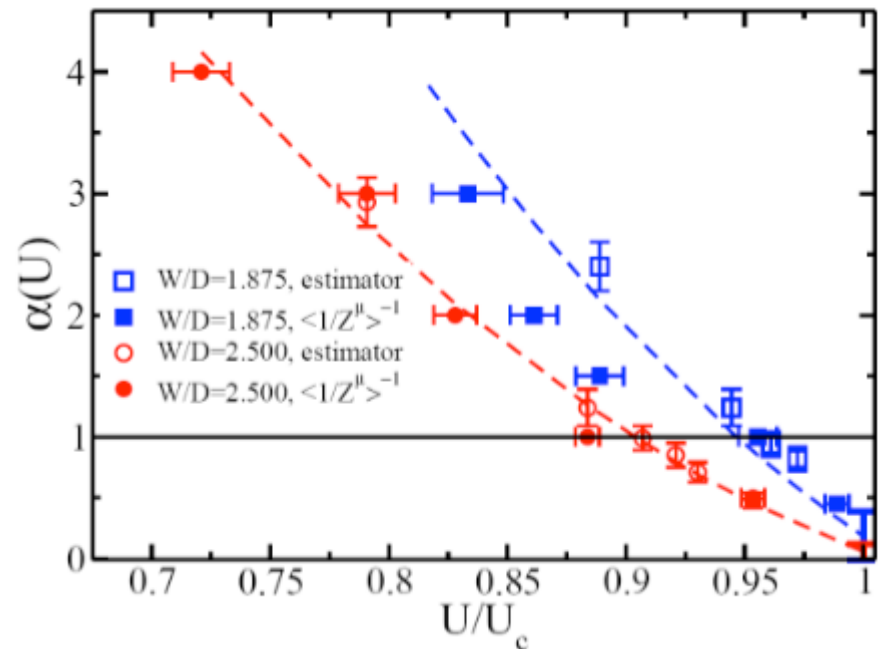
$$\chi_i(T) \sim \frac{1}{T + T_{Ki}} \Rightarrow \langle \chi(T) \rangle \sim \int dT_k \frac{T_K^{\alpha-1}}{T + T_K} \sim T^{\alpha-1}$$

Singular thermodynamic response

The exponent  $\alpha$  is a function of disorder and interaction strength.  $\alpha=1$  marks the onset of singular thermodynamics.

## Quantum Griffiths phase

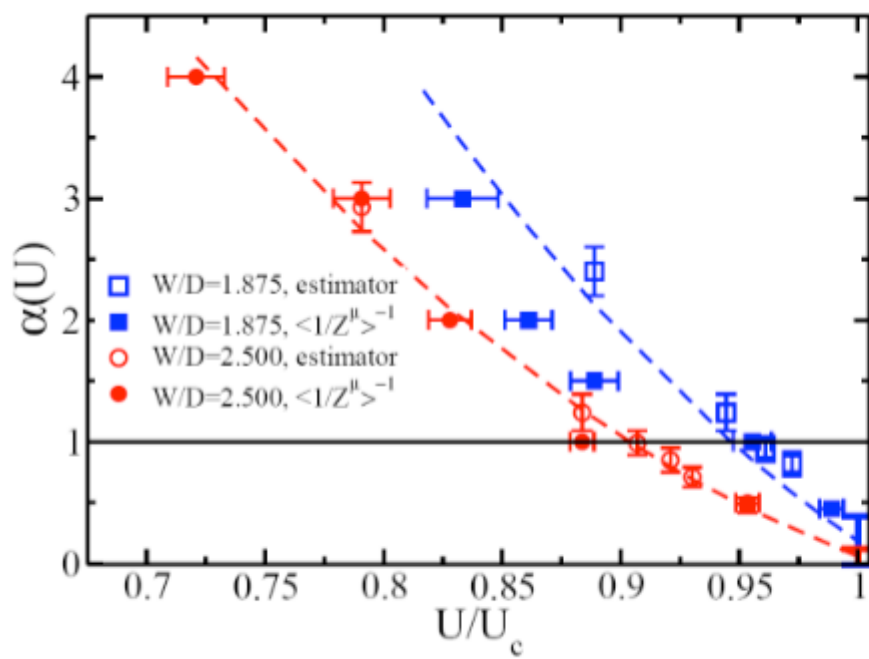
E. Miranda and V. D., Rep. Progr. Phys. **68**, 2337 (2005); T. Vojta, J. Phys. A **39**, R143 (2006)





## Infinite randomness at the MIT?

- Most **characterized** Quantum Griffiths phases are precursors of a critical point where the **effective disorder is infinite** (D. S. Fisher, PRL **69**, 534 (1992); PRB **51**, 6411 (1995); ....)



$$P(Z) \propto Z^{\alpha-1}$$

$$\alpha \rightarrow 0 \Rightarrow \Delta Z \rightarrow \infty$$

**Compatible** with  
infinite randomness  
fixed point scenario

**$1/\alpha$  – variance of  $\log(Z)$**

# “Size” of the rare events?

$$\chi_i \sim Z_i^{-1}$$

Replace the environment of given site outside square by uniform (DMFT-CPA) effective medium.

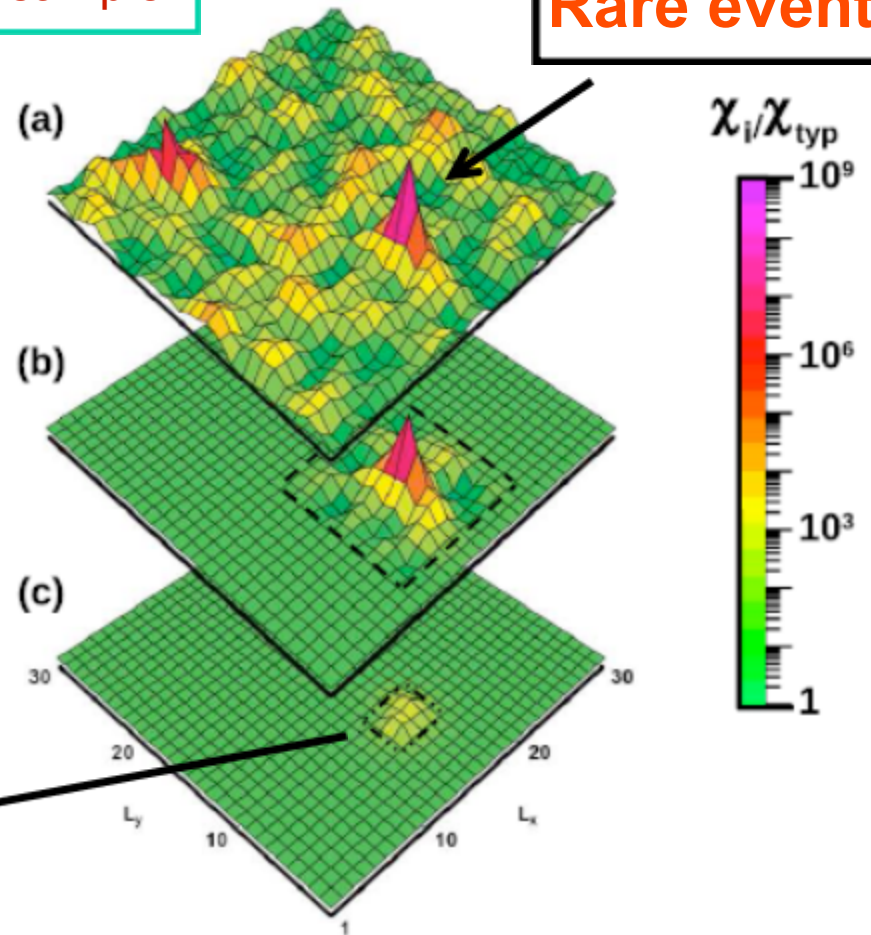
Reduce square size down to DMFT limit.

Rare events due to **rare regions** with weaker disorder

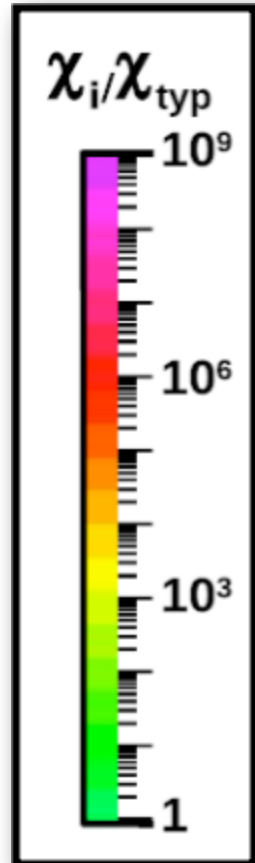
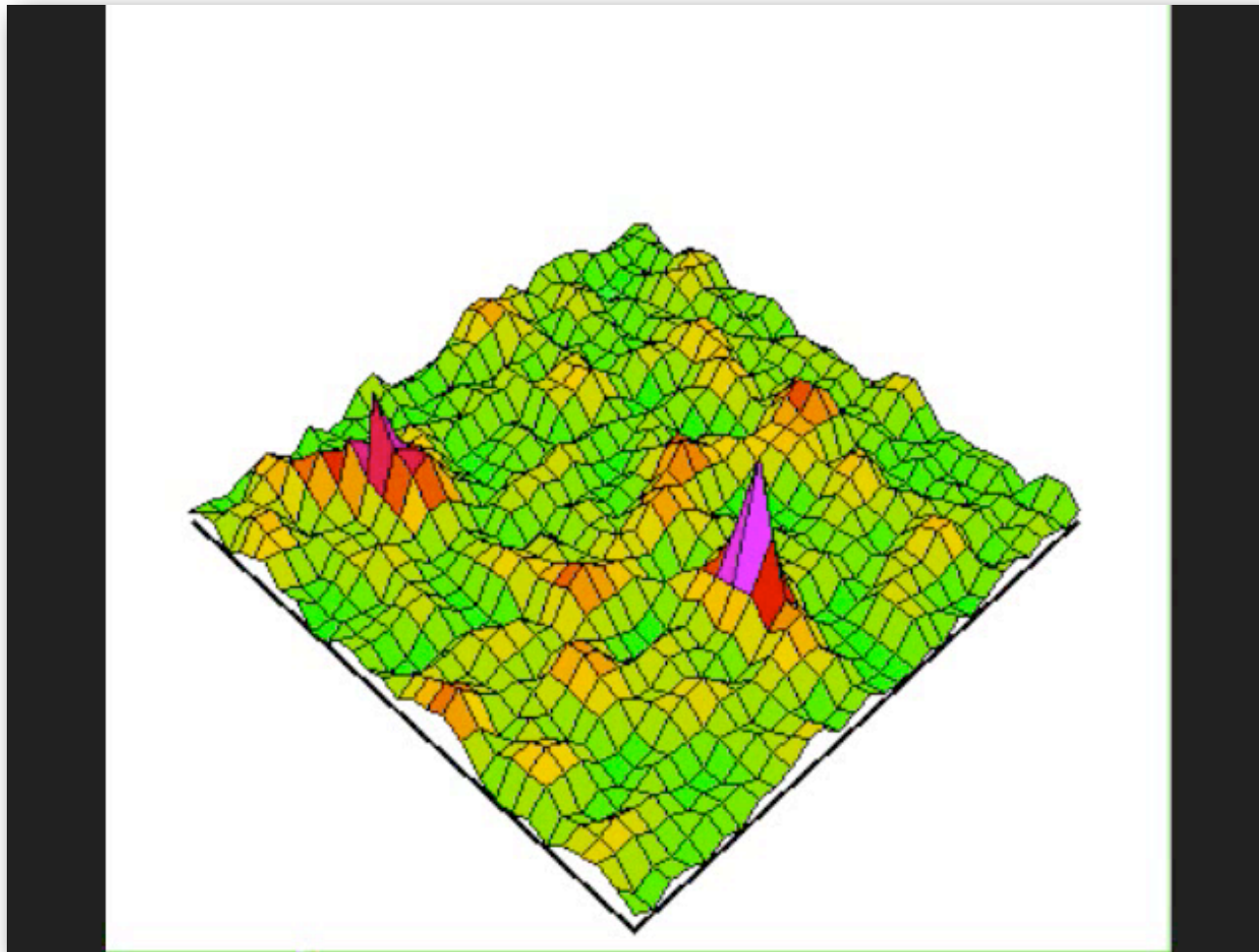
The rare event is **preserved** for a box of size  $l > 9$ : **rather smooth profile with a characteristic size.**

Typical sample

Rare event!



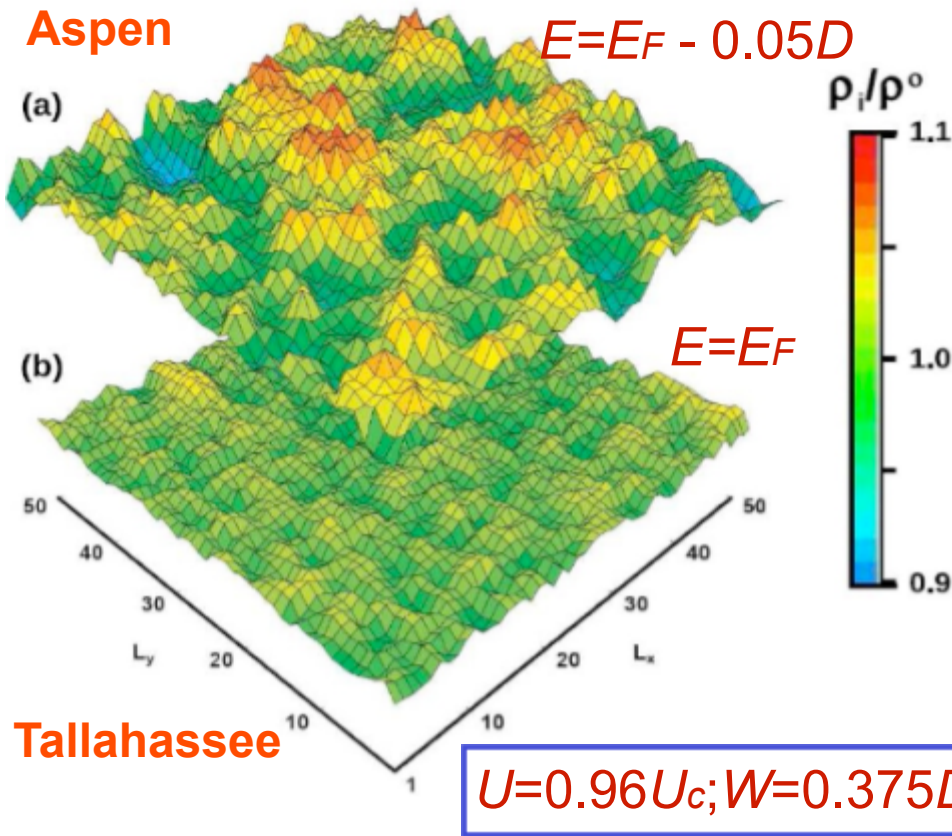
## “Size” of the rare events: a movie



**Killing the Mott droplet**

# Energy-resolved inhomogeneity!

- However, the effect is lost even slightly away from the Fermi energy:



$$v_i(\omega \neq 0) = \varepsilon_i + \Sigma_i(\omega \neq 0)$$

$$= v_i + \omega(1 - Z_i^{-1})$$

The strong disorder effects reflect the wide fluctuations of  $Z_i$

Similar to high- $T_c$  materials, as seen by STM

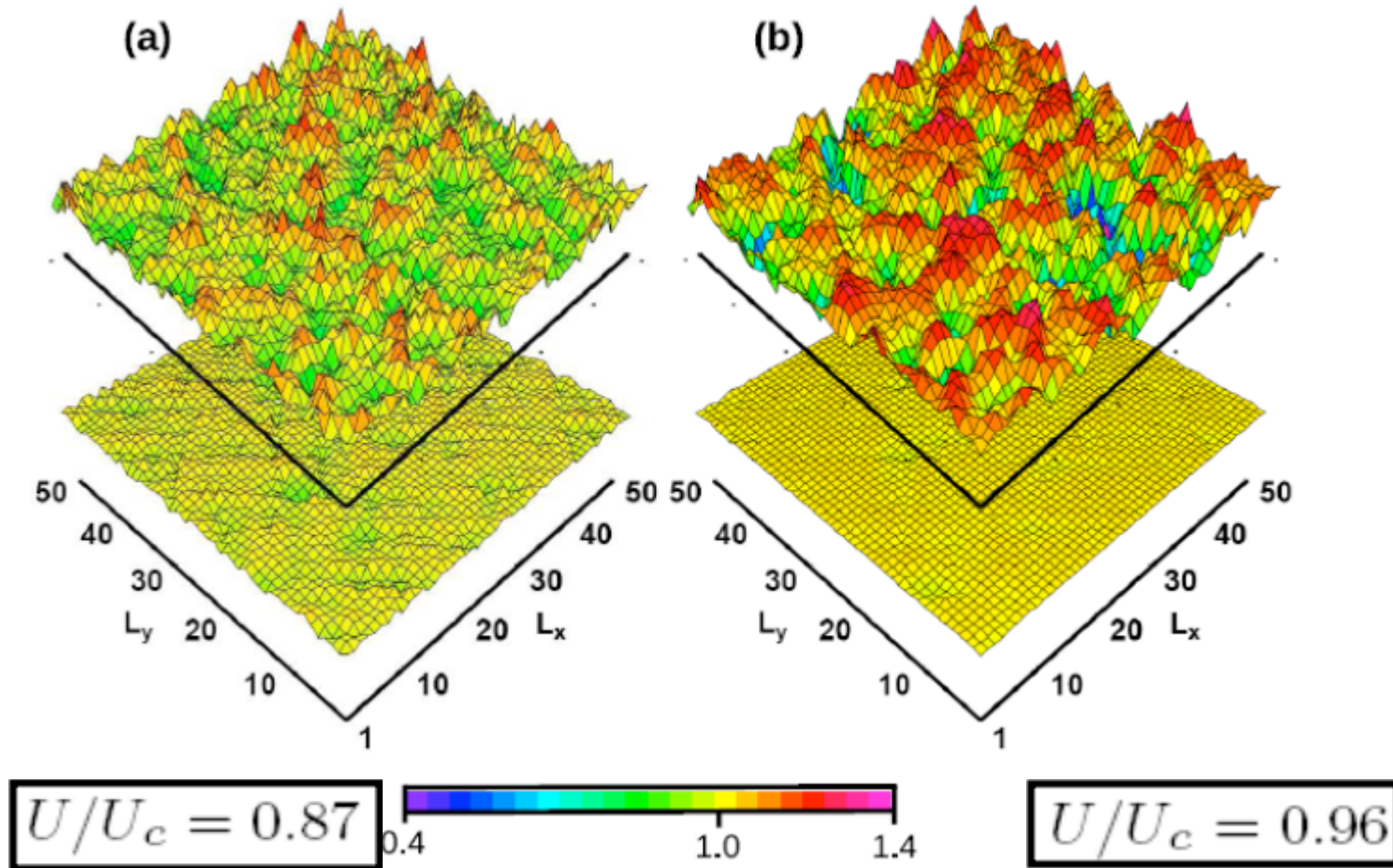
**Experiment: Seamus Davis (2005)**

**Theory: Garg, Trivedi, Randeria (2008))**

**Generic to the strongly correlated materials?**

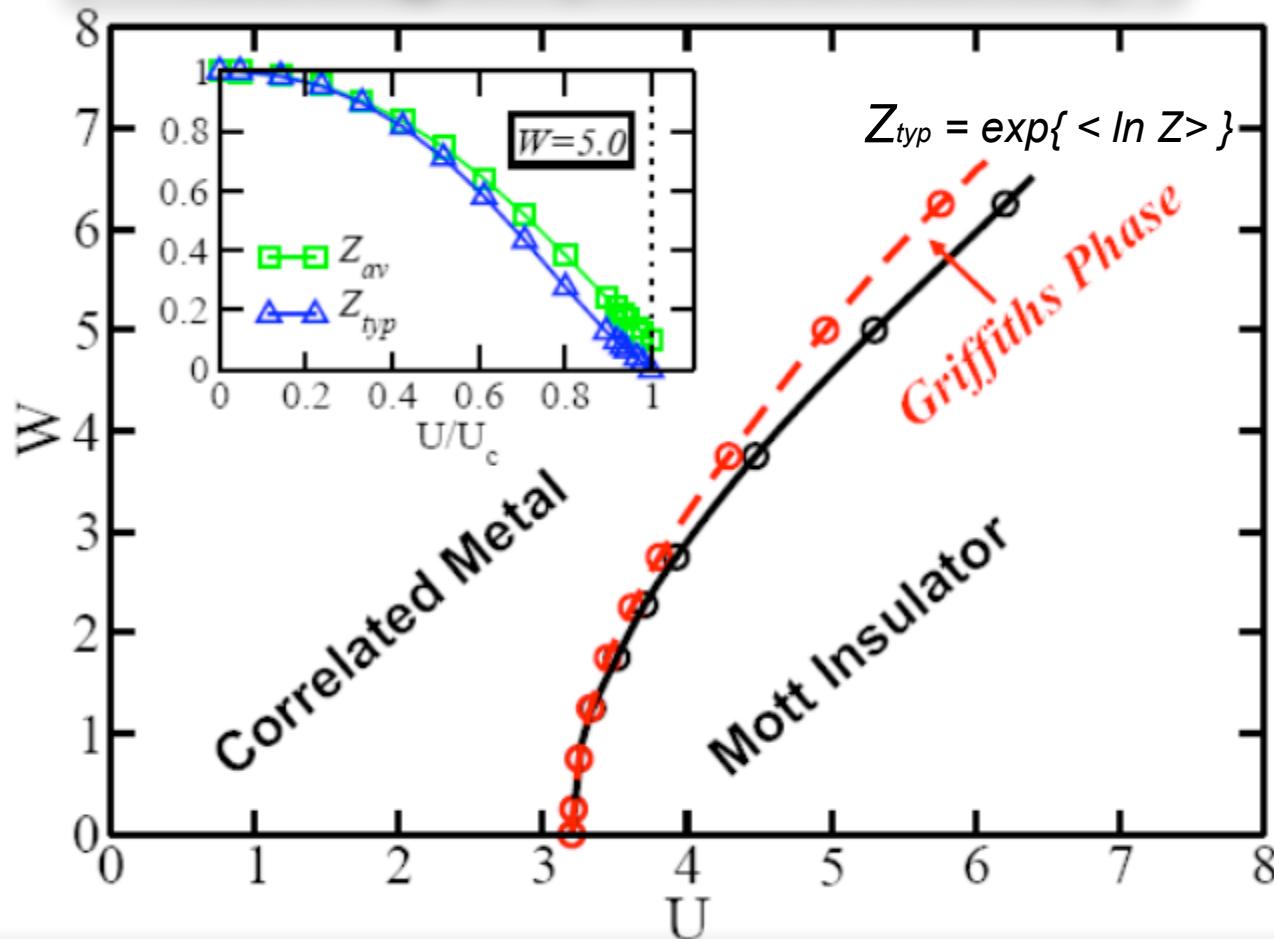


# Mottness-induced contrast



**Generic feature of all Mott systems, not only high  $T_c$  cuprates?!**

# Phase Diagram (moderate disorder)



**Fermi liquid**

**Non-Fermi liquid**

**IRFP (Mott-like?) insulator**

$\chi(T=0)=\chi_0$

$\chi(T=0)=\infty$

**MIT**

$\chi(T=0)=\infty$

**Disorder**

# NFL in Heavy Fermion Quasi-Crystals?

PRL 115, 036403 (2015)

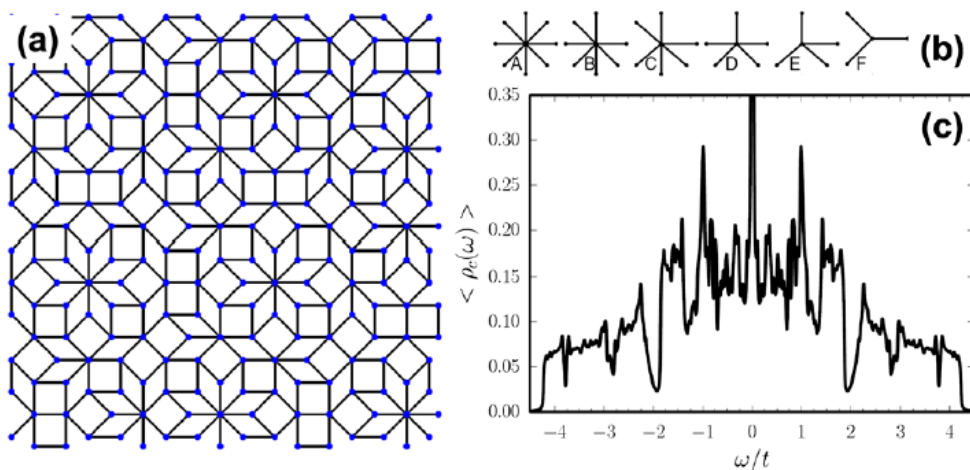
PHYSICAL REVIEW LETTERS

week ending  
17 JULY 2015

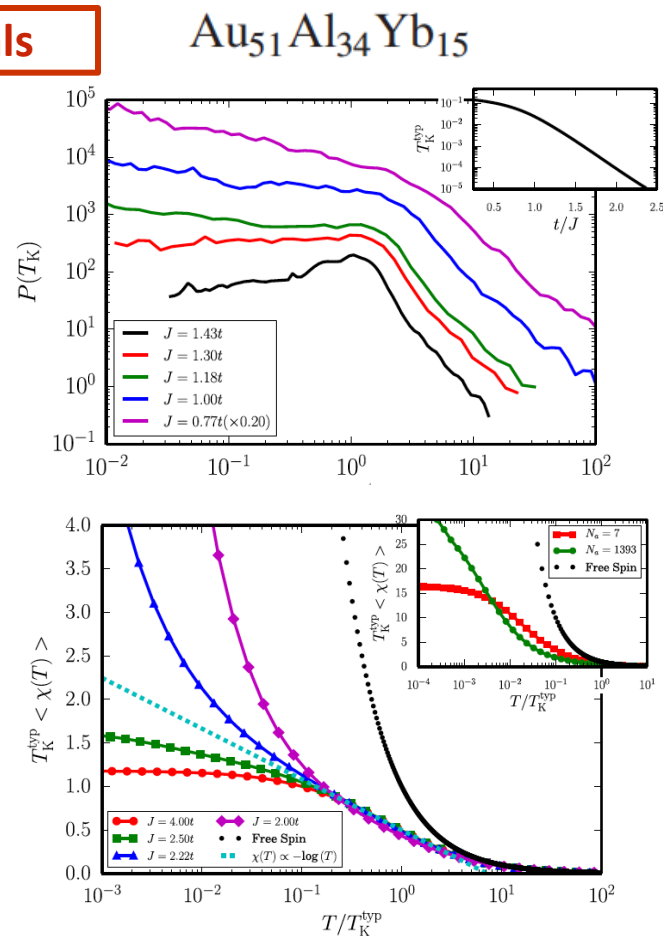
## Non-Fermi-Liquid Behavior in Metallic Quasicrystals with Local Magnetic Moments

Eric C. Andrade,<sup>1</sup> Anuradha Jagannathan,<sup>2</sup> Eduardo Miranda,<sup>3</sup> Matthias Vojta,<sup>4</sup> and Vladimir Dobrosavljević<sup>5</sup>

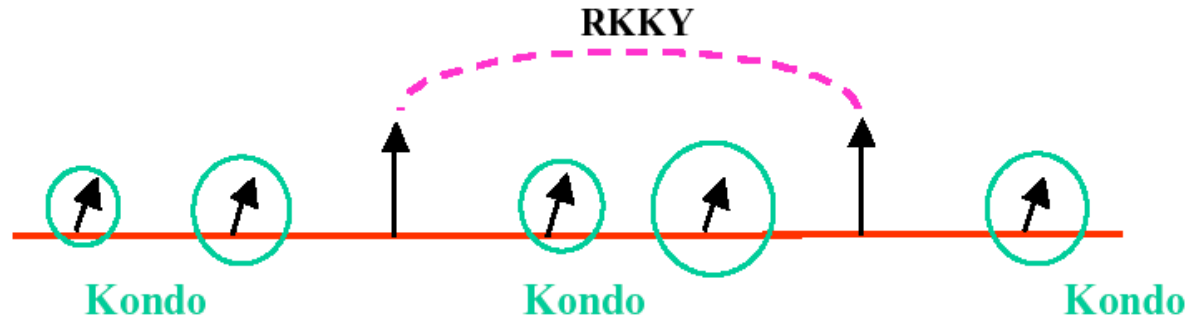
**New Materials: Rare-earth inter-metallics quasi-crystals**



**Aperiodic system - distribution of Kondo temps.  $T_K$   
Electronic Griffiths Phase, similar to disordered HF**



## Adding RKKY interaction: spin glass instability



- RKKY interactions between (distant) low- $T_K$  (unscreened) spins: oscillatory with distance  $\Rightarrow$  **random in magnitude and sign**
- Expect quantum **spin-glass** (SG) dynamics at low  $T$

**EDMFT** theory for RKKY interactions:

Bosonic bath: 
$$S_{RKKY} = g \int d\tau d\tau' \vec{\sigma}_f(\tau) \chi(\tau - \tau') \vec{\sigma}_f(\tau')$$

Self-consistency: 
$$\chi(\tau - \tau') = \overline{\langle \vec{\sigma}_f(\tau) \vec{\sigma}_f(\tau') \rangle}$$

Bose-Fermi (BF) Kondo model: **additional dissipation** from bosonic bath!



# Fractionalization and Two-Fluid Behavior

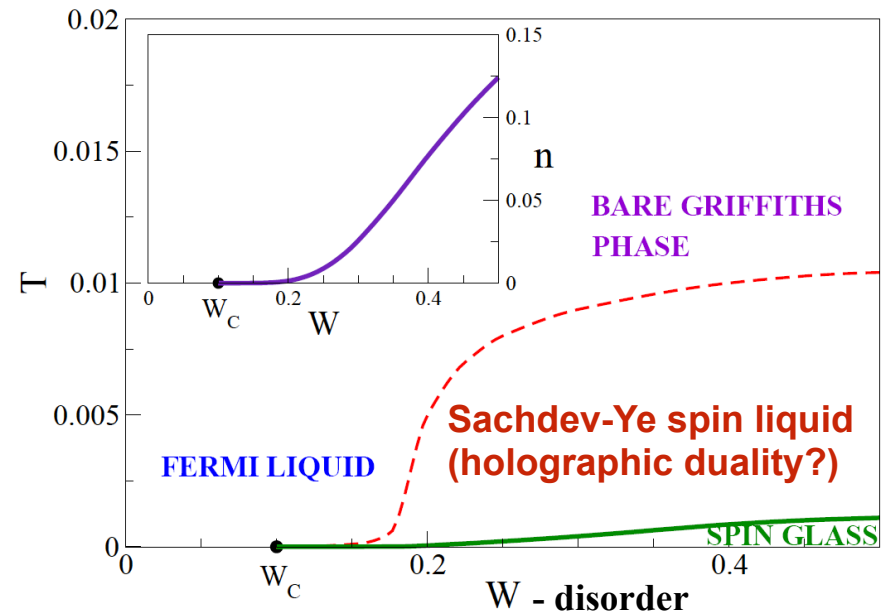
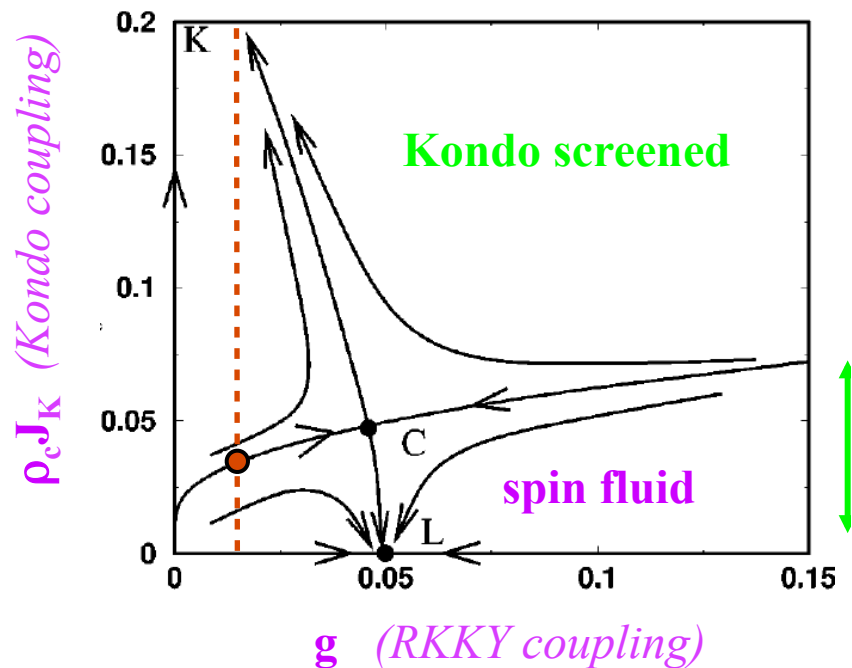
PRL 95, 167204 (2005)

PHYSICAL REVIEW LETTERS

week ending  
14 OCTOBER 2005

## Spin-Liquid Behavior in Electronic Griffiths Phases

D. Tanasković,<sup>1</sup> V. Dobrosavljević,<sup>1</sup> and E. Miranda<sup>2</sup>



# Analytical insight at weak disorder

PRL 104, 236401 (2010)

PHYSICAL REVIEW LETTERS

week ending  
11 JUNE 2010

## Quantum Ripples in Strongly Correlated Metals

E. C. Andrade,<sup>1,2</sup> E. Miranda,<sup>1</sup> and V. Dobrosavljević<sup>2</sup>



One impurity – **Friedel oscillations**

Interference: quantum corrections

Ballistic:  $\Delta\sigma \sim T$  (d=2)

Diffusive:  $\Delta\sigma \sim \log T$  (d=2)

(Aleiner, 2001)

Weak impurity - analytic (perturbative) solution, numerics - general

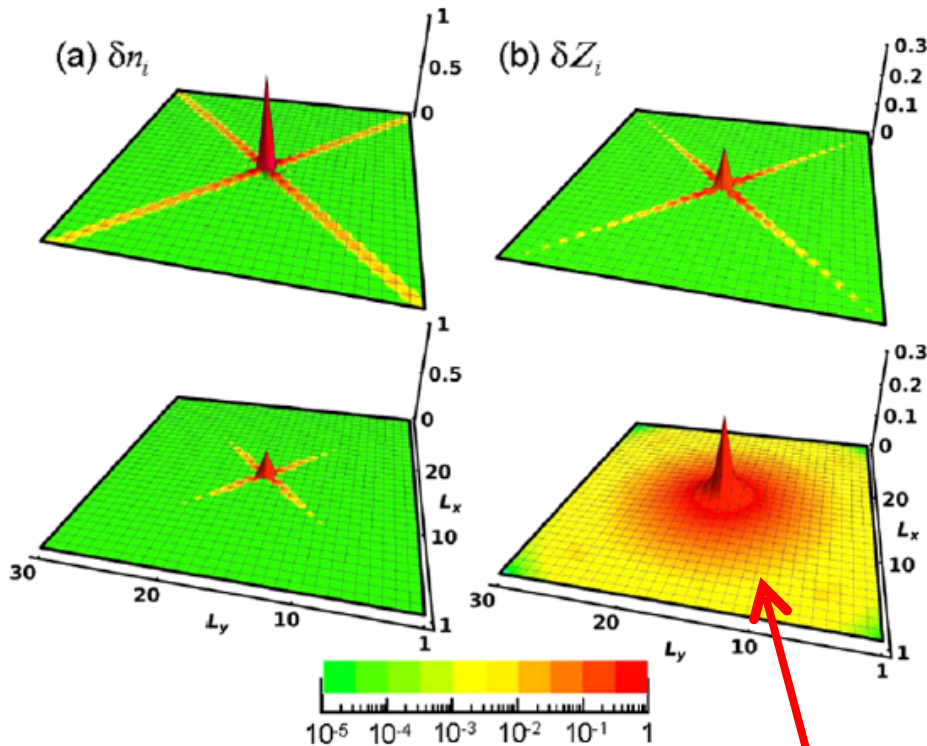
Reduce to standard Hartree-Fock results at small U

**Correlated regime and nonlocal terms??**

# Spatially correlated density fluctuations and Gutzwiller factors $Z_i$



## “Healing” length and Mott quantum criticality



$$\xi = (2z(1-u))^{-1/2}$$

“healing” length

“healing” length

$$\delta Z_i \sim \frac{1-u}{U_c^2} \left( \frac{\pi^{(1-d)/2}}{2^{(1+d)/2} \xi^{(d-3)/2}} e^{-r_{ij}/\xi} - 4(1-u)^3 [\Pi^{(0)}]_{ij}^{-1} \right) \varepsilon_j^2$$

“healing”

Lindard function

# Quantum corrections vs. inelastic scattering



## 2D ballistic regime

$$\tau_{\text{tr}}^{-1}(T) = \tau_0^{-1} A^2(u) \left\{ 1 + 2 \frac{T}{T_F} \alpha(u) w(T, \gamma(T)) + \eta \gamma(T) \right\}$$

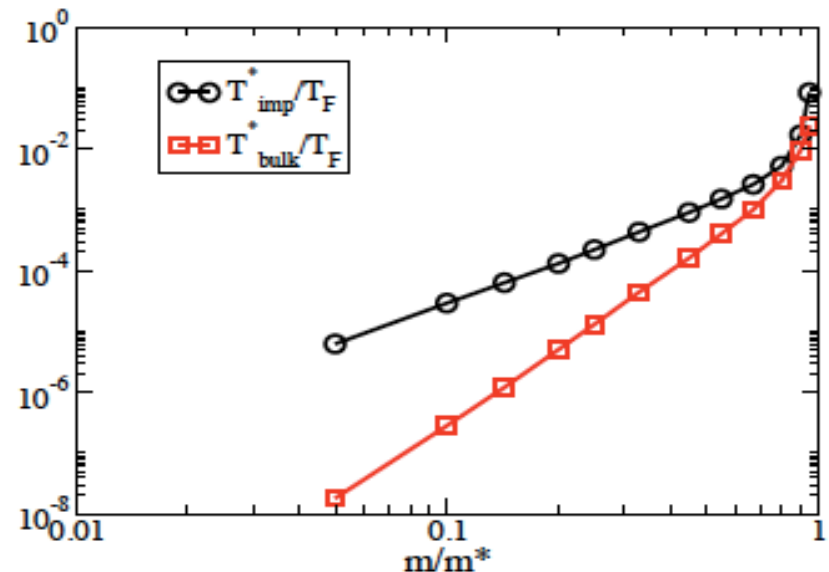
$$w(T, \gamma) = \int \frac{dx}{4} \text{Sech}^2 \left[ \frac{x}{2} \right] \text{Re} \left[ \ln \Gamma \left( \frac{1}{2} + \frac{\gamma(T)}{2\pi T} + i \frac{x}{2\pi} \right) \right] + \frac{1}{2} \ln(2\pi) + \frac{\gamma(T)}{2\pi T} \ln \left( \frac{T_F}{2\pi T} \right)$$

Linear T transport  
only at  $T < T^* \ll T_F$

## Inelastic (electron-electron) scattering

$$\gamma(T) = C \Lambda(u) T_F (T/T_F)^2$$

$$\Lambda(u) \sim (1-u)^{-2}$$



# Mottness-induced healing in strongly correlated superconductors

Shao Tang,<sup>1</sup> E. Miranda,<sup>2</sup> and V. Dobrosavljevic<sup>1</sup>

$$H = - \sum_{ij\sigma} t_{ij} c_{i\sigma}^\dagger c_{j\sigma} + J \sum_{ij} \mathbf{S}_i \cdot \mathbf{S}_j + \sum_i (\epsilon_i - \mu_0) n_i$$

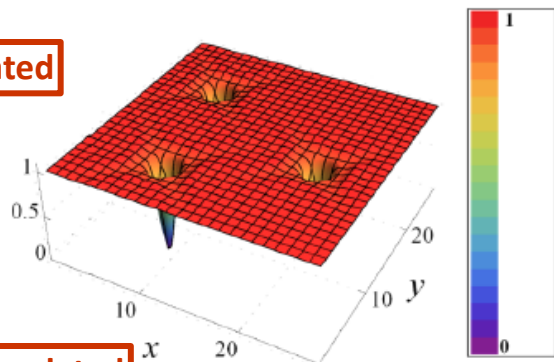
**Doped Mott insulator  
+ weak disorder**

D-wave pairing, large-N; Kotliar-Liu, 1992

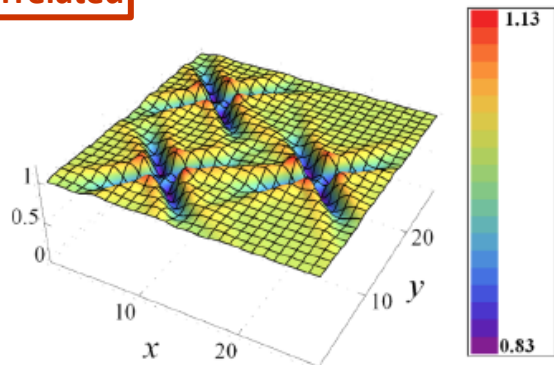
**gap function**

$$\Delta_{ij} = \langle f_{i\uparrow} f_{j\downarrow} - f_{i\downarrow} f_{j\uparrow} \rangle \quad \left\langle \frac{\delta \Delta_i}{\Delta_0} \frac{\delta \Delta_j}{\Delta_0} \right\rangle_{\text{disorder}} = f(\mathbf{r}_i - \mathbf{r}_j)$$

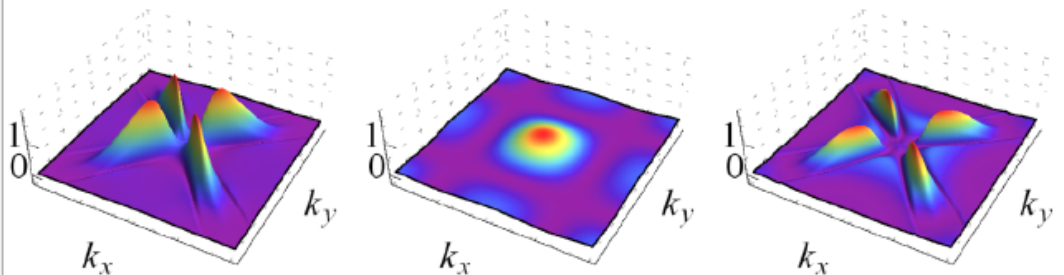
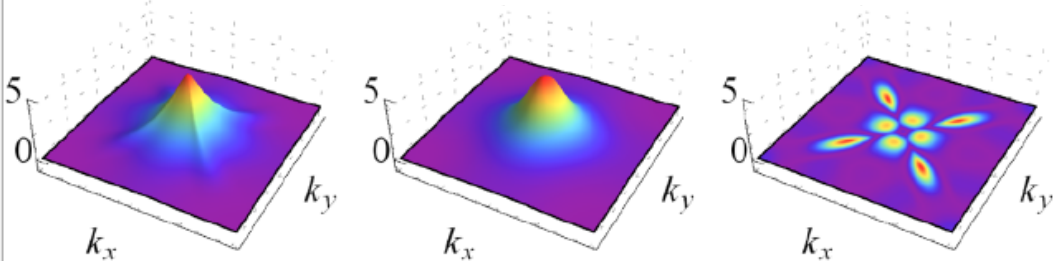
**correlated**



**un-correlated**



**correlator in momentum space**



# “Healing” vs. Abrikosov Gor’kov pair-breaking?

PHYSICAL REVIEW B 93, 195109 (2016)

**Strong correlations generically protect *d*-wave superconductivity against disorder**

Shao Tang,<sup>1</sup> V. Dobrosavljević,<sup>1</sup> and E. Miranda<sup>2</sup>

$$H = - \sum_{ij\sigma} t_{ij} c_{i\sigma}^\dagger c_{j\sigma} + J \sum_{ij} \mathbf{S}_i \cdot \mathbf{S}_j + \sum_i (\epsilon_i - \mu_0) n_i$$

D-wave pairing, large-N; Kotliar-Liu, 1992

**Doped Mott insulator  
+ weak disorder**

**BCS+AG: Even non-magnetic impurities -> strong pair-breaking for D-wave??**

T-matrix at weak disorder + strong correlations

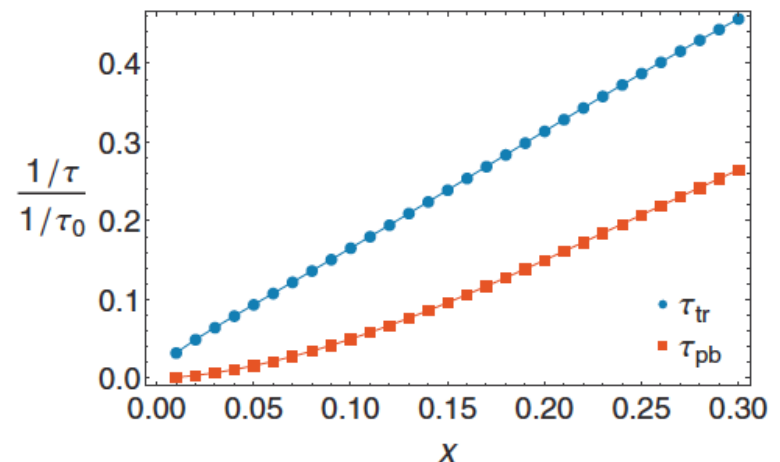
$$\ln \frac{T_{c0}}{T_c} = \psi \left( \frac{1}{2} + \frac{\alpha}{2} \right) - \psi \left( \frac{1}{2} \right)$$

$$\alpha \equiv 1/(2\pi T_c \tau_{pb})$$

$$\frac{1}{\tau_{pb}} = \frac{x^2 n m^*}{2\pi} \int_0^{2\pi} d\theta \, g \left[ \left| \sin \left( \frac{\theta}{2} \right) \right| \right] (1 - \cos 2\theta)$$

$$g(y) \equiv \frac{1}{\{\rho^* \lambda_0 k_F^2 y^2 g_L(y) + x[1 - 2\rho^* E_F g_L(y)]\}^2}$$

**Strong suppression of pair-breaking!**

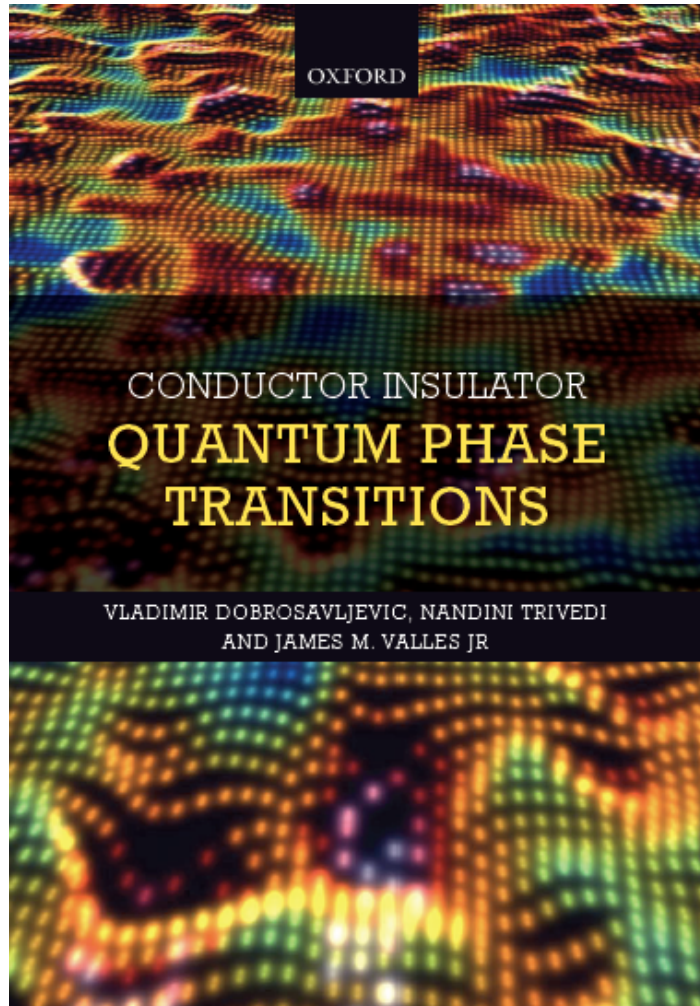


## Perspectives and challenges

- Significant disorder **renormalization** due to correlations
- Diffusion modes, “**interaction-localization**”?
- Disorder-induced **non-Fermi liquid** behavior
- Other **non-periodic** systems with correlations (HF quasi-crystals)?
- **IRFP** behavior, SDRG approaches?
- Inter-site (magnetic) correlations, **CDMFT**?
- Bosonic modes in “weak” FL, **fractionalization**, spin-glass **EDMFT**?
- Behavior **out of equilibrium (MBL)** with strong correlations and disorder?



To learn more:



<http://badmetals.magnet.fsu.edu>  
(just Google “Bad Metals”)

**Book:**

**Oxford University Press, June 2012**

**Already listed on Amazon.com**

**ISBN 9780199592593**